



# Evaluating a lightning parameterization based on cloud-top height for mesoscale numerical model simulations

J. Wong<sup>1</sup>, M. C. Barth<sup>2</sup>, and D. Noone<sup>1</sup>

<sup>1</sup>Cooperative Institute for Research in Environmental Sciences, Department of Atmospheric and Oceanic Sciences, University of Colorado, Boulder, CO 80309-0216, USA

<sup>2</sup>Atmospheric Chemistry Division, National Center for Atmospheric Research, Boulder, CO 80307-3000, USA

Correspondence to: J. Wong (john.wong@colorado.edu)

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**Abstract.** The Price and Rind lightning parameterization based on cloud-top height is a commonly used method for predicting flash rate in global chemistry models. As mesoscale simulations begin to implement flash rate predictions at resolutions that partially resolve convection, it is necessary to validate and understand the behavior of this method within such a regime. In this study, we tested the flash rate parameterization, intra-cloud/cloud-to-ground (IC : CG) partitioning parameterization, and the associated resolution dependency “calibration factor” by Price and Rind using the Weather Research and Forecasting (WRF) model running at 36 km, 12 km, and 4 km grid spacings within the continental United States. Our results show that while the integrated flash count is consistent with observations when model biases in convection are taken into account, an erroneous frequency distribution is simulated. When the spectral characteristics of lightning flash rate are a concern, we recommend the use of prescribed IC : CG values. In addition, using cloud-top from convective parameterization, the “calibration factor” is also shown to be insufficient in reconciling the resolution dependency at the tested grid spacing used in this study. We recommend scaling by areal ratio relative to a base-case grid spacing determined by convective core density.

simulated alongside of physics (e.g., Grell et al., 2005). Lightning-generated nitrogen oxides (LNO<sub>x</sub>) are predicted to be very efficient in accelerating the production of tropospheric ozone, which is identified as a significant greenhouse gas in the upper troposphere (Lacis et al., 1990; Kiehl et al., 1999). Cooper et al. (2007), showing that during the summertime North American Monsoon, lightning can contribute 25–30 ppbv of upper tropospheric ozone. Choi et al. (2009) has remarked on the increasing importance of LNO<sub>x</sub> in tropospheric ozone production as anthropogenic sources of NO<sub>x</sub> are being reduced in the United States. Furthermore, the inherent nonlinearity between NO<sub>x</sub> emission and commonly validated quantities such as radiative balances and ozone concentration makes it challenging to quantify the skill of a LNO<sub>x</sub> parameterization through proxy or total NO<sub>x</sub> measurements. Therefore, it is important to evaluate existing lightning parameterizations by directly validating flash rate predictions in order to more accurately interpret results from models that incorporate LNO<sub>x</sub> emission.

The most commonly used method for parameterizing lightning flash rate is perhaps that by Price and Rind (1992, 1993, 1994). It has been used by chemistry transport modeling studies such as GEOS-Chem (Hudman et al., 2007), MOZART-4 (Emmons et al., 2010), and CAM-Chem (Lamarque et al., 2012). Continental flash rates are related to the fifth-power of cloud-top height by Williams (1985) and Price and Rind (1992, hereafter PR92) through empirical evidences that are consistent with the theoretical scaling arguments of Vonnegut (1963). The partitioning between intracloud and cloud-to-ground flashes, or IC : CG ratio, is estimated with a fourth-order polynomial of cold cloud-depth,

## 1 Introduction

Over the last decade, predictions of lightning flash statistics in numerical weather and climate models have garnered increasing interests. One of the likely drivers is the advances in online chemistry models, wherein chemistry is

i.e., distance between freezing level and cloud-top, in Price and Rind (1993, hereafter PR93). Finally, the parameterization is generalized for different grid sizes with an extrapolated “calibration factor” in Price and Rind (1994, hereafter PR94).

Other bulk-scale or resolved-scale storm parameters may also be correlated with lightning flashes for the purpose of formulating alternative parameterization schemes. At parameterized convection scales, Grewe et al. (2001) used a mass flux approach derived by correlating observations of mass flux and cloud-top heights. Similarly, Allen and Pickering (2002) and Allen et al. (2010) implemented a parameterization of flash rate to the square of deep convective mass flux. Zhao et al. (2009) and Choi et al. (2005) based their flash rate predictions on both the deep convective mass flux and the convectively available potential energy (CAPE). Allen et al. (2012) used a flash rate predictions scheme based on the convective precipitation rate. Hansen et al. (2012) produced a lookup table for flash rate from convective precipitation and mixed phase layer depth by correlating data from observations. Appropriate for simulations with resolved convection, Petersen et al. (2005) gave a linear relation between flash rate and ice water path (IWP). Deierling and Petersen (2008) investigated a linear dependence of flash rate on updraft volume for  $T < 273\text{ K}$  and  $w > 5\text{ m s}^{-1}$ . Barthe et al. (2010) compared several of these methods, including PR92, through case studies, and showed that while the polynomial orders are lower in these formulations, the level of uncertainties may still be higher than PR92 due to a combination of errors from model biases in the parameters used, e.g., hydrometeors, and observational biases in the datasets used for constructing the relationships. Futyun and Del Genio (2007) arrived at a similar conclusion about the reduced reliability of precipitation-based approaches in global climate simulations for predicting lightning flash rate.

As a way to provide lightning hazard forecasts for the public in a qualitative manner, Yair et al. (2010) developed the lightning potential index (LPI) based on ice fractions and super-cooled liquid water mixing ratios between freezing level and  $-20^\circ\text{C}$ , and it has been shown to correlate well with observed flash rates in the Mediterranean. While the LPI does not directly produce a flash rate and no relationship was given to convert one from another, one of the many underlying assumptions is that charge buildup should be proportional to the fourth power of the relative velocities of the charging particles, strongly resembling the scaling arguments by Williams (1985). Similarly, Bright et al. (2005) introduced the Cloud Physics Thunder Parameter (CPTP) based on convective available potential energy (CAPE) and temperature at the equilibrium level (EL). Like LPI, CPTP is a qualitative index that does not translate directly to flash probability or flash count. Instead, a  $\text{CPTP} \geq 1$  is “considered favorable” for cloud electrification. McCaul et al. (2009) proposed two additional approaches to threat indices calculation based on vertical hydrometeor mass flux within the mixed-phase

region and column integral of frozen hydrometeor, which are currently operationally used in NOAA High-resolution Rapid Refresh (HRRR) model, a version of WRF.

The goals of this study are to evaluate the cloud-top height based parameterization (PR92, PR93, and PR94) across the bridging resolutions between those commonly used by global chemistry models ( $\Delta x \sim O(1^\circ)$ ) and cloud-resolving models ( $\Delta x < 5\text{ km}$ ), and report on statistics over time periods useful for studying upper tropospheric chemistry ( $O(\text{month})$ ) (Stevenson et al., 2006). It is, however, not the goal of this study to invalidate previous studies, but to draw attention to the need for careful implementation and validation of the use of these parameterizations. Here we report on experiments using PR92, PR93, and PR94 implemented into the Weather Research and Forecasting model (WRF; Skamarock et al., 2008), focusing on results from simulations performed at 36 km and 12 km grid-spacing. A simulation at 4 km grid spacing for 2 weeks in July and August 2006 is also analyzed to demonstrate how PR92 behaves transitioning from cloud-parameterized to cloud-permitting resolutions and provide insights on how or whether such transition can be done.

Similar studies have been performed for global models (e.g., Tost et al., 2007), but previous regional-scale modeling studies utilizing PR92 at comparable horizontal grid spacings have not provided evaluations of the lightning parameterization. Thus, there has been insufficient information to understand the behavior of PR92 in this regime. Even though these formulations were derived using near-instantaneous data at a cloud-permitting resolution (5 km), past applications often utilize temporally and spatially averaged cloud-top height outputs or proxy parameters. While the effects of spatial averaging is addressed by the PR94 scaling factor, effects of temporally averaging cloud-top heights are rarely addressed and may lead to significant underestimation due to the fifth-power sensitivity (Allen and Pickering, 2002). Addressing the potential issue of temporal averaging, instantaneous cloud-top heights and updraft velocities at each time step are leveraged. Comparisons are then performed for temporal, spatial, and spectral features.

The next section (Sect. 2) outlines the methods used in this study, which includes the formulation and overview of the parameterization (Sect. 2.1), relevant aspects of the model setup, practical considerations of implementing PR92 (Sect. 2.2), and the data used for validation (Sect. 2.3). Section 3 describes the model results and discusses the implications of various statistics from validation against observations of precipitation, flash rate, and IC : CG ratios. Section 4 discusses how the performance of PR92 transitions between different resolutions (Sect. 4.1) and between theoretically similar formulations (Sect. 4.2). Finally, Sect. 5 provides a summary of key results and cautionary remarks on specific aspects of the utilization of PR92, PR93, and PR94.

2 Methods

2.1 Parameterization overview

In PR92, a fifth-power relation between continental lightning flash rate ( $f_c$ ) and cloud-top height ( $z_{top}$ ) is established with observational data following the theoretical and empirical frameworks of Vonnegut (1963) and Williams (1985). Assuming a dipole structure with two equal but opposite charge volumes and a cloud aspect ratio of approximately one, it is first formulated, based on scaling arguments of Vonnegut (1963), that the flash rate would be proportional to maximum vertical updraft velocity ( $w_{max}$ ) and fourth-power of cloud dimension. Imposing a linear relation between  $w_{max}$  and cloud dimension, the flash rate relationship can be reduced to fifth power of  $z_{top}$  (Williams, 1985). It is empirically fit to radar and flash rate data from several measurements between 1960–1981 to give the continental equation (Price and Rind, 1992):

$$f_c(z_{top}) = 3.44 \times 10^{-5} z_{top}^{4.9} \tag{1}$$

PR92 also estimated that  $w_{max} = 1.49 z_{top}^{1.09}$  for continental clouds, thus allowing a second formulation based on maximum convective updraft:

$$f_c(w_{max}) = 5 \times 10^{-6} w_{max}^{4.54} \tag{2}$$

A separate formulation of second-order, instead of fifth-order, is also derived by Price and Rind (1992) for marine clouds, for which updraft velocity is observed to be significantly slower:

$$f_m(PR92)(z_{top}) = 6.2 \times 10^{-4} z_{top}^{1.73} \tag{3}$$

Taking into account effects from cloud condensation nuclei, Michalon et al. (1999) modified the marine equation to fifth-order:

$$f_m(M99)(z_{top}) = 6.57 \times 10^{-6} z_{top}^{4.9} \tag{4}$$

The practical viability of the continental relation was proven by Ushio et al. (2001) and Yoshida et al. (2009) through several case studies. However, Boccippio (2002) showed that the marine equations are formally inconsistent with Vonnegut (1963), and that the marine equations cannot be inverted to produce cloud tops within the range of cloud-top observations.

Price and Rind (1993) used the flash data from eleven states in the western United States, detected by wide-band magnetic direction finders, in combination with thunderstorm radar and radiosondes data to find a relation for the IC:CG ratio ( $Z$ ) from cold-cloud depth ( $d$ ), defined as the distance from freezing level to cloud-top.

$$Z = 0.021d^4 - 0.648d^3 + 7.49d^2 - 36.54d + 63.09 \tag{5}$$

Table 1. WRF simulations performed in this study.

Case #	dx (km)	dt (s)	Output	Duration
1	36	90	hourly	JJA 2006
2	36	90	hourly	JJA 2011
3	12	36	3-hourly	Jul 2011
4	4	12	hourly	25 Jul–7 Aug 2006

In Price and Rind (1994), a “calibration factor” ( $c$ ) for the resolution dependency of PR92 is introduced by regridding 5 km data between 1983 and 1990 from the International Satellite Cloud Climatology Project dataset (ISCCP; Rossow and Schiffler, 1991) to different horizontal grid sizes. The resulting equation is as follows:

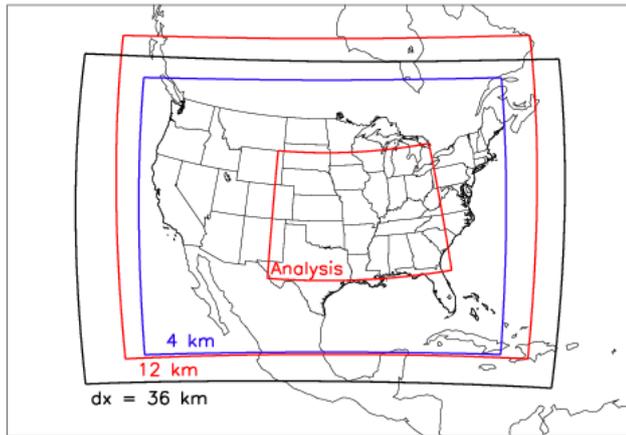
$$c = 0.97241 \exp(0.048203R) \tag{6}$$

where  $R$  is the grid area in squared degrees. Price and Rind (1994) claim that there is no dependence of  $c$  on latitude, longitude, or season. For the grid sizes used in this study, the values of  $c$  are 0.9774 for 36 km, 0.973 for 12 km, and 0.9725 for 4 km.

2.2 Model setup and implementation

Simulations in this study are performed using the Weather Research and Forecasting (WRF) model version 3.2.1 (Skamarock et al., 2008) over the contiguous United States (CONUS), including part of Mexico and Canada (Fig. 1). The simulations have slightly different model domains because the simulations were developed and performed for objectives independent of validating the lightning parameterization. Meteorology is initialized and continuously nudged (horizontal winds, temperature, water vapor) with the National Center for Environmental Prediction (NCEP) Global Forecasting System (GFS) final (FNL) gridded analysis at 6 h intervals (00:00 UTC, 06:00 UTC, 12:00 UTC, 18:00 UTC). Four simulations are performed (Table 1), two at 36 km grid spacing, one at 12 km grid spacing, and one at 4 km grid spacing. All cases use the same vertical coordinates with 51 sigma levels up to 10 hPa. The Grell–Devenyi ensemble convective parameterization (Grell and Devenyi, 2002) with Thompson et al. (2008) microphysics is used for the simulations where grid-spacing  $\Delta x > 10$  km, for which a convective parameterization is needed. The implementation of the GD scheme employed in this study consists of  $3 \times 3 \times 16 = 144$  ensemble members comprising of interactions between different dynamic control and static control/feedback closures. The maximum moist static energy (MSE) is then used as input with entrainment to calculate the level neutral buoyancy (LNB), or cloud top. For further information about the convective parameterization, readers are encouraged to refer to Grell (1993).

Since the simulations were designed independently, some physics options used are not consistent. The planetary



**Fig. 1.** Non-nested domains for WRF simulations and region for analysis.

boundary layer (PBL) parameterization is handled by the Yonsei University scheme (Hong et al., 2006) at 36 km and Mellon–Yamada–Janjic (MYJ) scheme (Janjić, 1994) at 12 km and 4 km. At 36 km, the surface layer physics option used is based on Monin–Obukhov similarity theory. The surface layer option used at 12 km and 4 km is also based on Monin–Obukhov theory but includes Zilitinkevich thermal roughness length.

While theoretically the scaling argument of Vonnegut (1963) does not distinguish between definitions of cloud-top height, the data used to derive the PR92 relation are radar reflectivity cloud-top heights at a certain reflectivity threshold. In the WRF implementation of Grell–Devenyi convective parameterization, the level of neutral buoyancy (LNB) is computed, with convective entrainment and detrainment accounted for within the calculation of cloud moist static energy, and readily available as a proxy for sub-grid cloud-top height. Thus, instead of 20 dBZ reflectivity cloud top,  $z_{\text{top}}$  is approximated by reducing LNB by 2 km, which will be shown to produce results within the range of the observed values. The choice of 2 km reduction is made independent of, but supported by, a recent study comparing different definitions of LNB and found the traditional “parcel” method definition of LNB over estimates the level of maximum detrainment by 3 km (Takahashi and Luo, 2012). Appendix A contains detailed discussions of the choice of 2 km cloud-top reduction and how it compares to offline computations of 20 dBZ cloud tops. Alternative methods for estimating the difference between the two heights can be formulated by directly taking into account their respective definitions. However, echoing Barthe et al. (2010), such addition of complexity increases the number of sources for uncertainty, especially in the context of parameterized convection. Similarly, using modeled cloud particle variables would also add an additional level of sensitivity due to sub-grid variability in hydrometeor mixing ratios. Therefore, reflectivity calcula-

tions are only performed in the 4 km simulation and only for the purpose of redistributing lightning flashes horizontally as described below.

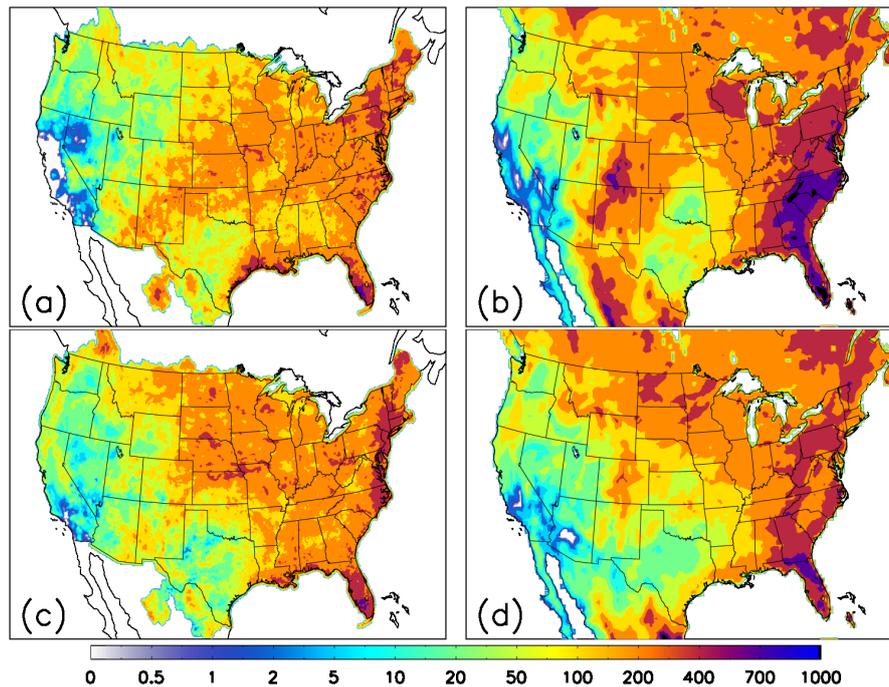
For case 4 (Table 1), convection is explicitly simulated with a modified Lin et al. (1983) microphysics scheme. Since no convective parameterization is used, the resolved maximum vertical velocities ( $w_{\text{max}}$ ) within the convective core are utilized (Barth et al., 2012), and Eq. (2) is used instead of Eq. (1) for estimating flash rate. In addition, since a single storm may often cover multiple model grids, flashes are redistributed to within regions with a minimum reflectivity of 20 dBZ calculated using hydrometeor (rain, snow, graupel) information that is now better constrained at 4 km. The IC : CG ratio is prescribed using a coarse version of the Boccippio et al. (2001) 1995–1999 climatological mean, which was computed using data from the Optical Transient Detector (OTD; Christian et al., 1996) and the National Lightning Detection Network (NLDN; Cummins and Murphy, 2009). Because PR92 developed Eq. (2) based on data at 5 km resolution, no resolution scaling is done to this simulation. Because this particular simulation was driven by the meteorology of its own WRF outer domains, it is restarted “cold” on 2 August to be consistent with the outer domain meteorology.

Most of the implementations used in these simulations are arguably “untuned” and not scaled to climatology or observations by any additional tuning factors, with the exceptions of the 2 km cloud-top height reduction used in the cases with parameterized convection and the prescribed climatological IC : CG ratios in case 4. Therefore, the correctness and predictiveness of the flash rate parameterization are not guaranteed at the time of the simulation given the lack of supporting validations of PR92 at the tested grid spacings. However, without feedback to the meteorology (except in case 4) and providing sufficient linearity in the biases of flash prediction, offline comparisons should reveal any tuning requirements for operational and research uses.

### 2.3 Data description

While desirable, event-by-event analysis would be technically challenging because the simulation may not produce the same strength, timing, and location of each convective event. Furthermore, an event-by-event analysis is unnecessary in the context of a mesoscale upper tropospheric chemistry study, of which the meaningful timescales often average biases from many individual events. Therefore, a large area where thunderstorms commonly occur is selected. The “analysis domain,” defined as 30°–45°N, 80°–105°W (Fig. 1), is used for time series and statistical comparisons.

The predicted lightning properties depend strongly on how the model simulates convection. Thus, in Sect. 3.1, WRF simulated precipitation is compared against National Weather Service (NWS) precipitation products to evaluate the model’s skill in representing convective strengths. The data are collected from radars and rain gauges and improved



**Fig. 2.** Spatial distribution of 2006 and 2011 JJA total precipitation in millimeters. (a) and (c) are NWS precipitation degraded to 12 km resolution. (b) and (d) are 36 km WRF-simulated total precipitation over the same periods with data above water surfaces masked out.

upon using a Multi-sensor Precipitation Estimator (MPE). Manual post-analyses are then performed by forecasters to identify systematic errors ([http://www.srh.noaa.gov/abrfc/?n=pcpn\\_methods](http://www.srh.noaa.gov/abrfc/?n=pcpn_methods)). The final data products used here are mosaic CONUS precipitation maps from 12 River Forecast Centers (RFCs) during JJA 2006 and 2011. The data are gridded into 4 km resolution and are available as 24 h totals over a hydrological day beginning and ending at 12:00 UTC.

The simulated CG flash counts, computed online as predicted total flashes  $\times$  predicted CG fraction, are compared against data from the Vaisala US National Lightning Detection Network (NLDN; Cummins and Murphy, 2009). The network provides continuous multiyear CONUS and Canada coverage of  $> 90\%$  of all CG flashes with ongoing network-wide upgrades (Orville et al., 2002, 2010). The median location accuracy is 250 m, which is well within the resolutions employed in this study. Multiple strokes are aggregated into a single flash if they are within 1 s and no more than 10 km apart. Weak positive flashes with  $< 15$  kA have been filtered from all data. Finally, the flash data are binned into hourly flash counts for each model grid cell for comparison against model output.

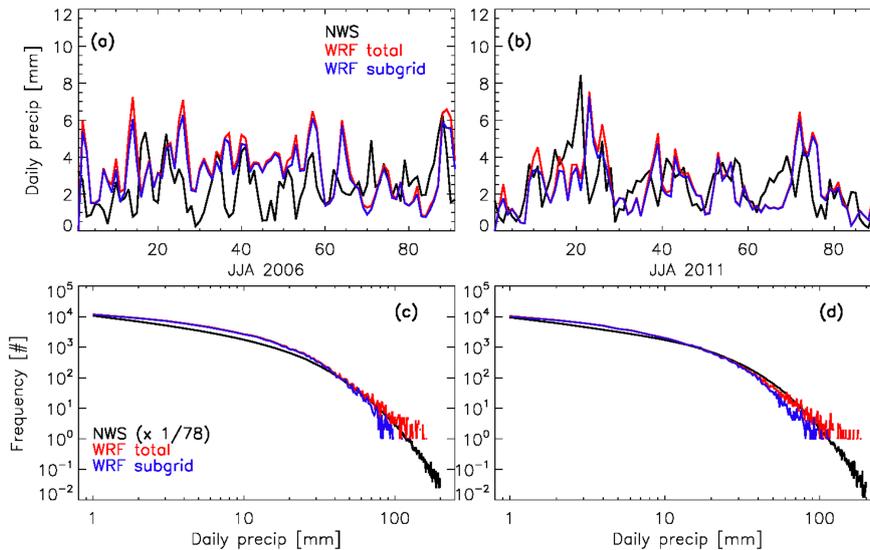
Data from Earth Networks Total Lightning Network (ENTLN), previously WeatherBug Total Lightning Network (WTLN), are used to validate the model-produced IC:CG ratios. ENTLN employs a wide-band system that operates between 1 Hz to 12 MHz (Liu and Heckman, 2011). The theoretical detection efficiency (DE) for CG flashes across

CONUS is 90–99 %, while the IC DE falls between 50–95 % (50–85 % within the analysis domain). Since the mappings of the corresponding DEs are not available with the data, 95 % and 65 % are used for CG and IC DEs, respectively, in analyses for which a prescribed DE is required. To address the concern of the impact of this simplification, the range of possible flash counts, IC:CG ratios, and biases will be provided when appropriate within the discussion in Sect. 3 to place bounds on the uncertainty. Due to the limited deployment duration of the network, only the IC:CG ratios during JJA 2011 within the analysis domain (see Fig. 1) are estimated and compared. For consistency with the comparisons against NLDN CG flash counts, the stroke aggregation criteria used here are 10 km and 1 s as done by NLDN, instead of the 10 km and 700 ms window typically used by Earth Networks to generate flash statistics.

### 3 Results and discussions

#### 3.1 Precipitation

While lightning does not directly depend on precipitation, they are both the results of the same processes that promote ice–graupel collisions. Further, precipitation is observed robustly and continuously, thus giving us a high quality measurement for validating model results. On the other hand, while convective mass flux may produce a more consistent correlation with lightning, the lack of well-controlled direct



**Fig. 3.** Time series and frequency distributions for JJA 2006 and 2011 area-averaged daily precipitation within the analysis region (see Fig. 1). Distributions for NWS are scaled by the ratios between total grid counts in WRF at 36 km and total grid counts in NWS within the analysis boundaries ( $\sim 1/78$ ). WRF subgrid is the portion of precipitation from subgrid cumulus parameterization. Only grid points with more than 1 mm of precipitation are included.

observations and the large uncertainty in model calculations make it an inferior proxy for convective strength in this context.

Figure 2 shows the total precipitation during JJA 2006 and 2011 over the CONUS as simulated at 36 km grid spacings by WRF and observed by NWS. The gradients across the CONUS for both years are well captured by the model, but WRF has a high bias for 2006. WRF also simulates up to an order of magnitude more precipitation for coastal regions for both years but primarily for 2006. The time series for mean daily area-averaged precipitation and frequency distributions for JJA 2006 within the analysis domain (Fig. 3a and c) also reveal a median model bias of 37%. In particular, WRF predicted more than twice the precipitation between late-June and mid-July in 2006. In contrast, the median bias for 2011 is 4.9% with almost equal occurrence of over- and under-predictions. The model frequency distribution for both years also closely track those observed (Fig. 3c and d) except at the high end of the distribution where limits of model grid resolution induces significant noise.

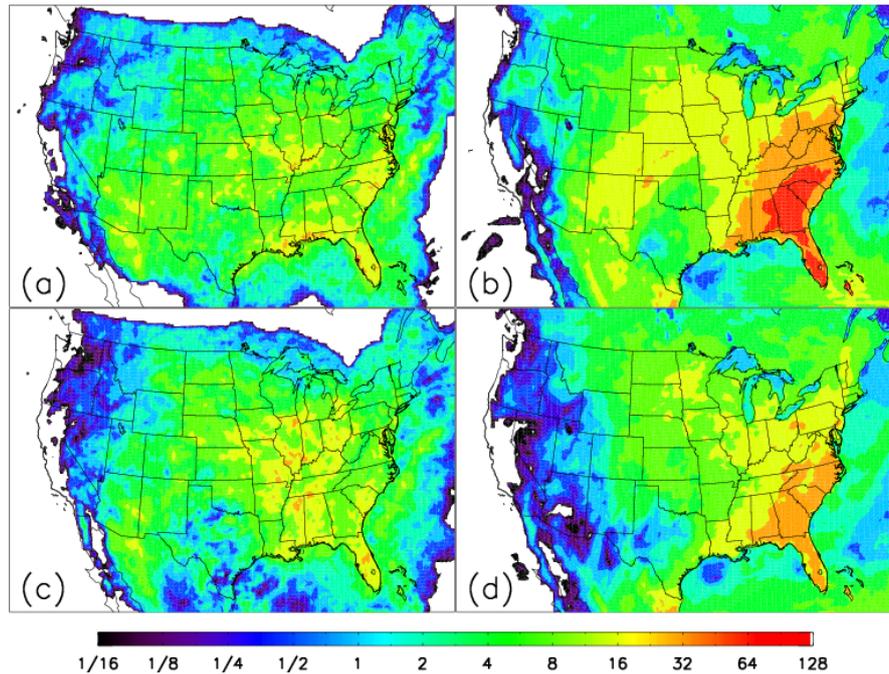
The simulated daily precipitation at 12 km is higher than the NWS observed precipitation by 24% during July 2011. However, an anomalously strong diurnal cycle is simulated at 12 km grid spacing that is not present in the 36 km simulation. Comparing the area-averaged 12 km nocturnal precipitation over the entire analysis domain to that of 36 km output, nocturnal precipitation at 12 km is too low but the daytime precipitation is too high. One-day simulations were performed to evaluate the impact from the differences in model physics, but there remained significant unidentified

discrepancies between the precipitation amount in the two runs that cannot be explained by horizontal resolution differences alone; thus, it is concluded that there is no value in re-doing the entire simulation. The identified causes for the differences between the two simulations are, in decreasing order for magnitude of influence, initial conditions for soil temperature and soil moisture, differences in planetary boundary layer scheme (Sect. 2.2), and the land surface model option. Such difference in diurnal behavior in the simulations is expected to have significant impact on how the lightning parameterization is evaluated, but the full impact can be minimized through incorporation of precipitation into the analysis.

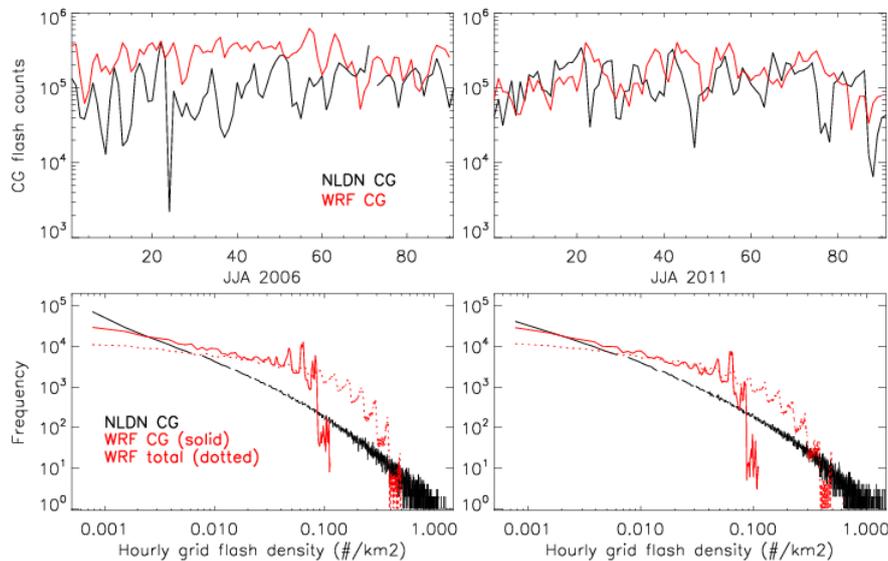
Large scale meteorology and moisture inputs are unlikely the causes for these biases due to nudging. In 2006, biases mostly occurred in the low-to-mid end of the distribution (i.e., light precipitation events, Fig. 3c), indicating that the problem lies in parameterizing subgrid events. Despite the differences in the eastern United States, convection over the central United States is reasonably represented. Finally, the goal of this study is not to evaluate the convective parameterization nor specific model setup, but rather to evaluate the performance characteristics of the lightning parameterization when implemented into a regional model with all the expected (and unexpected) defects.

### 3.2 CG flash rate

Figure 4 shows the CONUS CG flash density (units in number per  $\text{km}^2$  per year). WRF is consistently higher along the East Coast for 2006 where positive bias is also observed



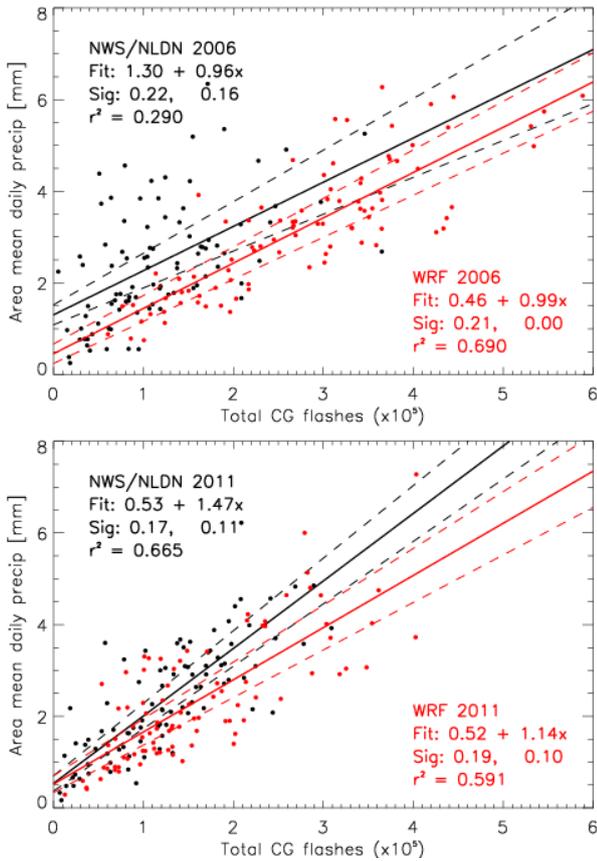
**Fig. 4.** Total CG flashes in number per km<sup>2</sup> per full-year during JJA 2006 (first row) and 2011 (second row). First column (a and c) shows the NLDN observed density gridded to WRF 36 model grid, and second column (b and d) shows the modeled flash density output by WRF at 36 km.



**Fig. 5.** Comparisons of time series and frequency distributions between NLDN CG flash counts (black) and WRF predicted CG flash counts (red) at 36 km within the analysis domain defined in Fig. 1. Total flash counts predicted by WRF are shown as dotted red lines.

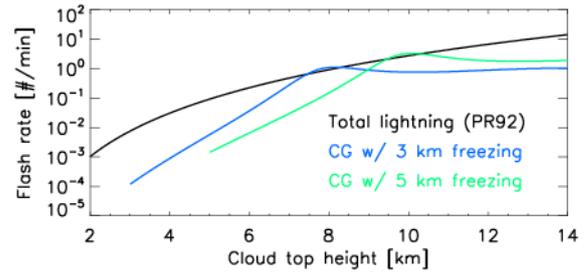
in the modeled precipitation, which is used as a proxy for quantifying the comparison of simulated convective strength against observations. Similarly, both flash rate and precipitation are over-predicted in the Colorado and New Mexico region for 2006. On the other hand, the low precipitation bias

in Arizona simulated by WRF for 2011 is coincident with a severe low bias in the same region for the CG flash density. Otherwise, flash densities are within the order of magnitude of those observed for regions where simulated precipitation is consistent with NWS observations.



**Fig. 6.** Total CG flashes (#) versus area-mean daily precipitation (mm) within the analysis domain (Fig. 1). Solid line is the least-square linear fit and dashed lines are  $\pm 1\sigma$  for both the constant terms and first-order coefficients. WRF is simulated at 36 km.

The over-prediction of CG flash density along the East Coast in 2006 dominates the regional mean and produces significantly high biases compared to 2011. Figure 5a and 5b show the time series of the total number of ground flashes predicted by WRF and observed by NLDN within the analysis region (Fig. 1). The median daily CG bias is 140% for 2006 and only 13% for 2011. It should be noted that these values were obtained by sampling all 3 months. Sampling one month would produce varying results. For instance, while JJA 2011 produces an overall median bias of 13%, July 2011 alone produces about twice as much lightning as observed but is offset by under-predictions in June and August. Because the lightning detection efficiency of NLDN varies spatially, the CG bias can vary over ranges of 116–154% for 2006 and 1.7–20% for 2011. The differences between the median biases for the two summers can be attributed largely to the differences in the total precipitation biases, as illustrated in the previous section (Sect. 3.1), for which 2006 is 37% too high while 2011 is only 5% higher than observations.

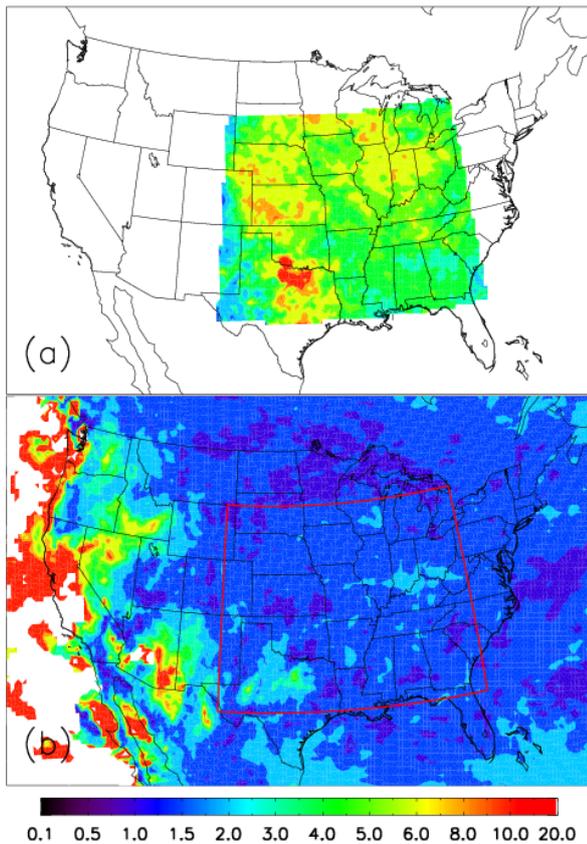


**Fig. 7.** Total lightning and CG flash rates computed using PR92 and PR93 for various cloud-top heights and freezing levels, demonstrating the source of spectral cut-off in Fig. 6.

To take into account the bias in the simulated convective strength, area-averaged daily precipitation is correlated with total CG flash count. While the relation is likely nonlinear, the area averages over the analysis domain are roughly linear in both WRF-simulated and observed data (Fig. 6). The slopes for the 2006 data are statistically the same, there is a constant positive bias for model produced flash counts over observed values. In contrast, 2011 results are close for small values but modeled and observed values diverge for more intense events. Such inconsistency between years demonstrates the potential for strong inter-annual variability in the correlation between flash rate and precipitation.

Figure 5c and d show the frequency distributions of the hourly grid flash density. From the spectra, it is apparent that the over-prediction observed in the time series occurs between flash densities of 0.003 to 0.1 CG flashes  $\text{km}^{-2} \text{h}^{-1}$ . However, the abrupt cutoff beyond  $\sim 0.11$  in both 2006 and 2011 modeled distribution indicates that PR92 fails to replicate the observed distribution. The occurrence of this cutoff can be explained by the local maximum when combining the PR92 total flash rate parameterization and PR93 IC : CG ratio parameterization (Fig. 7). Together, the predicted CG flash rate is capped at a certain limit, depending on the freezing level regardless of the cloud-top height. In addition, the total flash rate is also under-predicted for high flash rate events (dotted red lines in the figures), thus contributing to the truncated model frequency distribution.

An initial comparison of the model results against the Lightning Imaging Sensor (LIS) total flash data also shows a high bias ( $\sim 2.3\times$ ) in overall flash count but an underestimation of the high flash count events, similar to the results demonstrated by the NLDN comparison. However, comparisons against the LIS data for this study have a low confidence level because of the relatively short time period simulated and the many uncertainties, such as variable detection efficiency and shifting diurnal sampling bias of LIS data, associated with the analysis.



**Fig. 8.** IC : CG bulk ratios for JJA 2011 as (a) observed by ENTLN and (b) predicted by WRF at 36 km grid spacing using PR93. The ENTLN detection efficiency used here is 0.65 for IC and 0.95 for CG.

### 3.3 IC : CG ratio

The JJA 2011 IC : CG bulk ratios ( $\equiv \sum_t IC(x, t) / \sum_t CG(x, t)$ ) are calculated within the analysis domain (Fig. 8a) using constant detection efficiencies of 95 % and 65 % for CG and IC flashes, respectively. While WRF produced a median IC : CG ratio of 1.74 within the region, ENTLN observed a median of 5.24 with a possible range of 3.80 to 7.17 due to the spatial variability in both IC and CG DEs. Considering the ambiguity in the choice of cloud-top definition described in Sect. 2.2, a possible solution to increase the IC : CG ratio computed using Eq. (5), thus achieving better comparison against observations, is by eliminating the cloud-top height reduction, an option that maintains the conceptual interpretation of the parameterization but has the potential of offsetting the bias. For consistency, the cloud-top height used in the total lightning parameterization needs to be un-adjusted as well.

To learn whether reasonable lightning flash rates and IC : CG ratios can be estimated by using just the level of neutral buoyancy (LNB), an offline calculation is made of the

daily flash counts with the cloud-top height adjustment eliminated. The offline calculation is performed using instantaneous, hourly model output of LNBs and temperatures (for determining freezing levels). While the offline calculation is able to replicate almost precisely the online flash count prediction, which causes both time series to appear overlapping in Fig. 9, the CG flash rate frequency distribution is severely degraded because of vertical discretization of cloud tops to model levels and lowered temporal resolution to hourly outputs. When LNB is used for the cloud-top height (with no adjustment), the prediction of both CG and total lightning flash rates increase, as expected. The CG median bias over ENTNLN increases from 44 – 51 % to 158 – 172 %, and the total lightning median negative bias of 53 – 25 % becomes a positive bias of 23 – 95 % for the aforementioned range of DEs. Furthermore, even though the frequency distribution of total lightning is closer to the observed distribution, the CG distribution still experiences the truncation as described in Sect. 3.2.

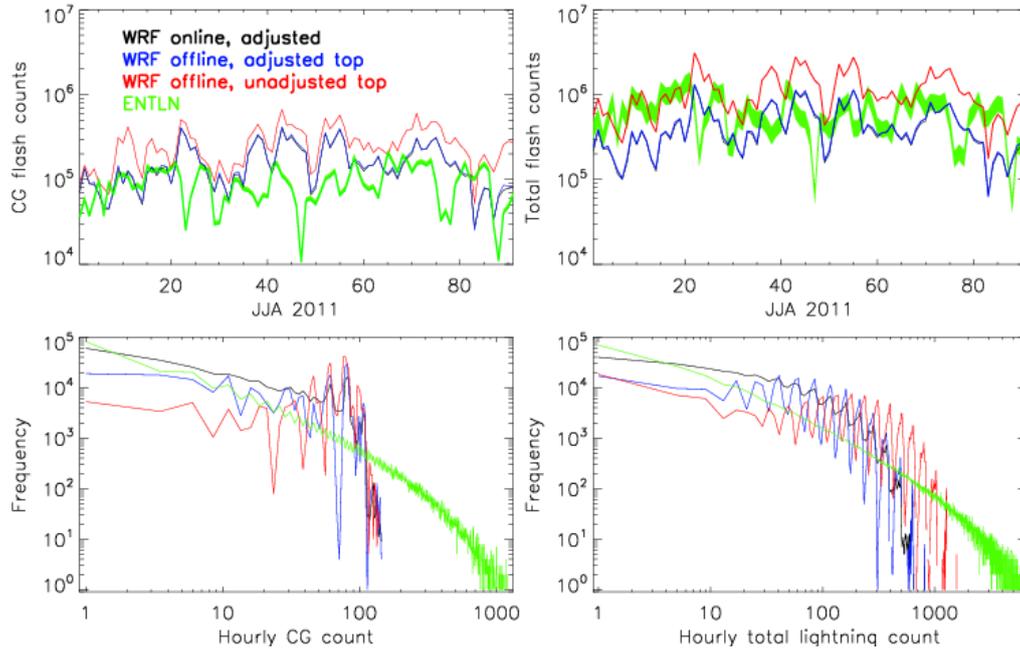
## 4 Resolution dependency

A goal of this study is to evaluate the applicability of the PR92 parameterization to resolutions between fully parameterized and partially resolved convection. Thus, it is useful to evaluate how the parameterization behaves as the grid size changes. To test the behavior of the PR94 calibration factor, a 12 km simulation for July 2011 is used. As grid sizes are reduced to allow convective parameterization to be turned off, the transition to  $w_{max}$  based formulation of PR92 (Eq. 2) is tested with a 4 km simulation between 25 July–7 August 2006. The domains for these simulations are shown in Fig. 1. Together, the results from these simulations will provide insights and recommendations on how to achieve resolution-awareness or independence while using PR92.

### 4.1 Sensitivity to grid size

At 12 km, the resolution dependency factor or “calibration factor” ( $c$ ) from Price and Rind (1994) is 0.56 % smaller than that applied to 36 km. However, comparison against the 36 km simulation and observations shows that there is a factor of  $\sim 10$  high bias. While there are differences in the statistics of convective strengths between the two simulations, as quantified by precipitation in Sect. 3.1, they are too minor to fully reconcile the large bias at 12 km. Therefore, an areal ratio scaling factor ( $1/9 = 12^2/36^2$ ) is applied offline to partially reconcile the differences on top of  $c(\sim 1)$ , which was applied online.

There are two reasons why the use of areal scaling instead of PR94 is justified in this study. The first reason pertains to why PR94 failed while it has been shown to work in GCMs. The PR94 calibration factor was derived from area-averaged cloud-top heights for progressively larger grid sizes from the



**Fig. 9.** Comparison of WRF predicted lightning flash counts generated online and offline with and without  $-2$  km cloud-top height adjustments against ENTLN CG and total flash counts. Thicknesses of the ENTLN bands in the time series are computed using the minimum and maximum theoretical IC and CG detection efficiencies within the analysis domain. Noisiness of offline calculated distributions are associated with using hourly outputs only rather than accumulating flashes at every model time step. It should be noted that online (black) and offline (blue) WRF outputs with adjusted top appear coincident in the time series, but are evidently different from the frequency distribution.

original ISCCP 5 km resolution to  $8^\circ \times 10^\circ$ . On the contrary, the LNBs from the convective parameterization are expected to change only slightly with grid resolution as long as the environmental parameters such as buoyancy remain similar.

The second reason addresses why the areal ratio is expected to work for regional scales. PR92 produces flash counts in unit of number of flashes per storm; thus, when approaching almost convection-resolving resolutions, where major storm size is comparable to grid size, the appropriate scaling should be done according to the expected number of convective cores per grid. Since  $\Delta x = 36$  km gives a reasonable flash rate compared to observations over 3 months, we assumed one storm per grid at this resolution and scaled the flash counts from  $\Delta x = 12$  km as an areal ratio. However, the base case resolution may spatially vary because of the dynamics controlling the minimum-permitted distance between convective cores. At coarse grid sizes, the area covered by the convective storm systems may only be a fraction of the grid cell area. Thus, the area scaling ratio may not be applicable when changing from base-case grid spacing (with approximately one storm per grid) to much coarser grid sizes. A possible solution is to include a cloud fraction estimate as part of the scaling factor between grid sizes.

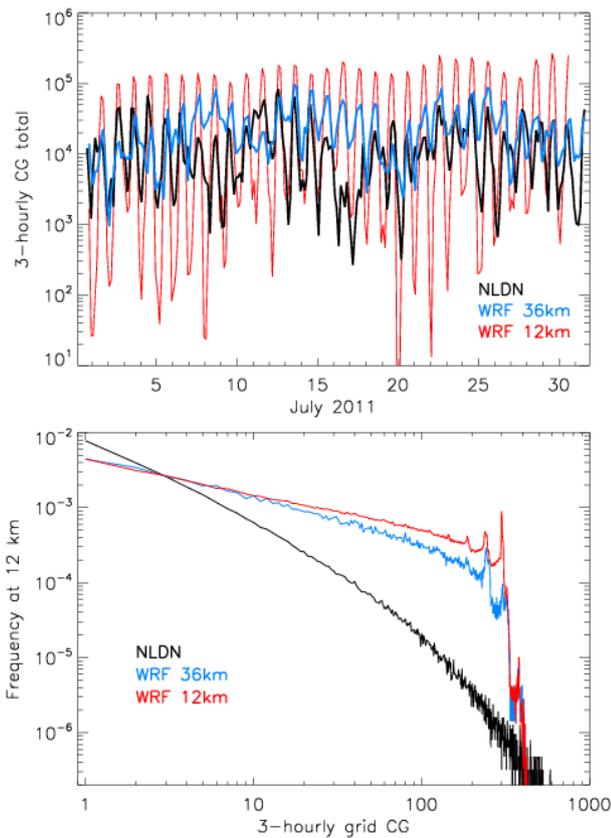
After scaling by  $1/9$ , WRF at 12 km predicts a median of 40% more 3-hourly lightning flashes than observed by NLDN (Fig. 10). This is to be compared with 36 km, which

predicted double the 3-hourly lightning for the same period. Simulating an anomalously strong diurnal cycle in precipitation, the 12 km flash count also shows a much more prevalent diurnal variation, associated with the poor simulation of the diurnal cycle of precipitation as previously noted. Much of the over-prediction is compensated by the negative biases in the nocturnal flash rates in the final statistics. Despite the differences in diurnal skill, the parameterization was able to produce the same drop-off in grid frequency distribution beyond 200 flashes per grid per 3 h, for which the primary cause is discussed in Sect. 3.2.

#### 4.2 Sensitivity to formulation

Comparing the 36 km simulation to the 4 km simulation provides insight into how the predicted flash density changes between resolutions using  $f(z_{\text{top}})$  for parameterized convection and  $f(w_{\text{max}})$  for resolved convective systems. This is an important factor to be considered if flash rate predictions are to be included in nested simulations or models permitting non-uniform grid-spacings such as Model for Prediction Across Scales-Atmosphere (MPAS-A; Skamarock et al., 2012).

The area-averaged daily precipitation predicted by the 4 km WRF-Chem simulation is 70% too high prior to 2 August 2006 and only 7.5% too high after 2 August. On 2 August, the 4 km WRF simulation was re-initialized (with no

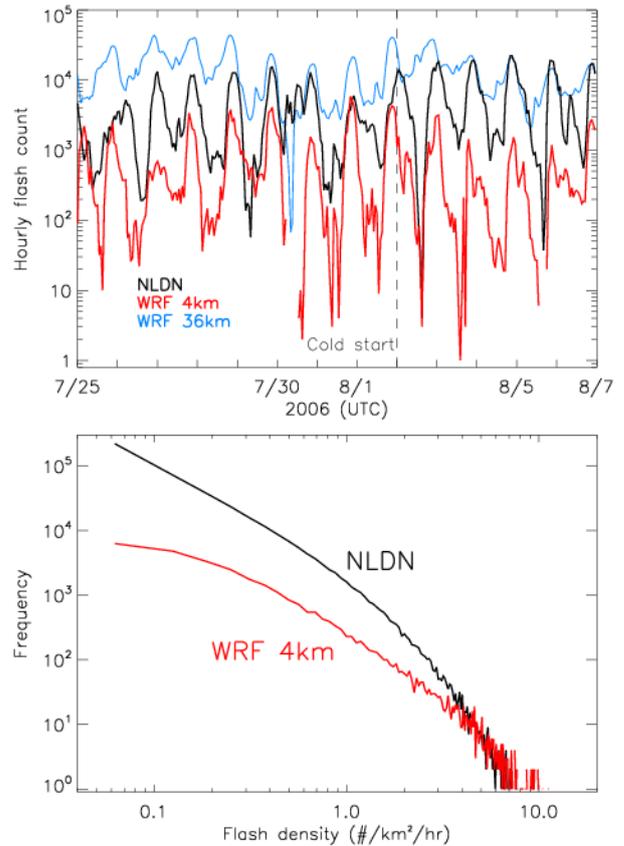


**Fig. 10.** Time series and frequency distributions of 3-hourly CG flash counts compared to NLDN at gridded to 12 km. The WRF 36 km distribution is adjusted by  $\times 9$  to account for the grid per area difference. The choice of computing the distributions for flash rate per grid as opposed to flash density is to demonstrate the consistency of the spectral drop-off at different resolutions.

clouds) to be consistent with the re-initializations of the outer domain WRF simulations that drove this 4 km simulation described in Barth et al. (2012). The flash rate predicted by the 4 km simulation follows the precipitation trend. A 26 % decrease in flash rate occurs between the period before 2 August and the period afterwards.

While the 36 km simulation over-predicted lightning flash rate for this period (25 July–7 August 2006), the 4 km simulation under-predicted the flash rate, exhibiting a  $-83\%$  bias relative to the NLDN flash counts prior to the cold-start and a  $-95\%$  bias after (Fig. 11). Similar underestimation of the  $w_{\max}$  formulation has been noted for both tropical (Hector storm near Darwin, Australia) and US continental storms (Cummings et al., 2013). These results indicate that it is important to evaluate the flash rate parameterizations with observations. It is insufficient to use high resolution model results as “truth” for coarse resolution simulations.

Despite the low bias in flash rate prediction, the 4 km WRF-Chem simulation matches the observed distribution of



**Fig. 11.** Time series and frequency distributions of hourly CG flash counts within the analysis domain as observed by NLDN and simulated by WRF at 4 km grid spacing.

flashes for high flash rate events and placed the burden of underestimation on the low-end of the distribution, which causes the distribution to appear flatter than observed. Since we are using a constant IC : CG ratio based on Boccippio et al. (2001) climatology instead of the PR93 parameterization, the erroneous drop-off in the CG flash rate distribution found in the other cases using PR93 is not present. Such improvement in spectral characteristics suggests that constant climatological IC : CG ratios may be a reasonable if not superior alternative to PR93.

### 5 Conclusions

We have implemented the WRF-Chem model parameterizations for lightning flash rate using prescribed IC : CG ratios and the associated resolution dependency by Price and Rind (1992, 1993, 1994), which are based on cloud-top height. In our implementation, the cloud-top height is estimated by the level of neutral buoyancy (LNB), adjusted by  $-2\text{ km}$  to reconcile the difference between LNB and radar reflectivity cloud top. No additional tunings and changes to the parameterizations are done. The modeled precipitation and

lightning flash rate are evaluated for the simulations with 36 km, 12 km, and 4 km grid spacings over CONUS for JJA 2006 and 2011.

The first result is that, after a 2 km reduction, the use of LNB as a proxy for cloud-top simulated at 36 km grid spacing produces CG flash rates at the same order of magnitude as NLDN observations. For models using other convective parameterizations, alternative choices of cloud-top proxies may be available and thus the appropriate methods of cloud-top adjustment should be determined on a case-by-case basis. Taking into account model biases in convection, as quantified by precipitation, the precipitation–lightning relation from the model and observations are statistically indistinguishable. While there is up to a factor of 2.4 median bias in the flash counts from the 2006 36 km simulation, it is accompanied by a 37 % over-prediction in precipitation. In contrast, the 2011 36 km simulation has a precipitation bias of 5 %, which leads to a 13 % over-prediction in flash counts. For the 12 km simulation the lightning flash rate bias is linked to the anomalously strong diurnal cycle simulated for convection, indicated by precipitation. Such bias in the simulated convection may be caused by a number of other model components.

Second, despite the correct CG count, it is shown that PR92 is incapable of producing the correct frequency distribution flash CG flashes, which are truncated at a much lower flash density than observed. The most likely cause is the function form of combining the PR92 total flash rate parameterization and the PR93 IC : CG ratio parameterization, which produces an upper-limit in the permitted maximum CG flash rate. This brings into question the validity of PR93 in contexts where spectra characteristics are a concern. It is recommended that using constant bulk ratios such as the climatology presented in Boccippio et al. (2001) or one derived from total lightning measurements may produce equal, if not better, spectra. Considering that the observed JJA 2011 IC : CG ratio also displays significant departure from the Boccippio et al. (2001) climatology for certain areas, it would be useful to revisit the subject of IC : CG climatology in future studies, taking advantage of the advances in continuous wide-area lightning detections over the past two decades.

Third, due to the use of LNB from the convective parameterization instead of area-averaged cloud-top heights, using the PR94 factor to adjust for different horizontal resolutions is not applicable and an areal ratio factor should be used instead to reconcile the resolution dependencies. Since the 36 km base cases produced relatively satisfying results, the 12 km simulation is scaled with an areal ratio relative to 36 km. However, it may be argued that the outcome from the 36 km simulations is only a corollary of the probability of having exactly one convective core within a single model grid. Therefore, other choices of “base case” grid spacings near 36 km may also produce similar results for CONUS, specifically within the analysis domain (Fig. 1), and other ar-

eas with different storm density may require a different base-case resolution for scaling. On the other hand, area ratios may not be appropriate at coarser resolution as convective core number density is highly non-uniform.

Finally, at 4 km, we used a theoretically similar formulation of PR92 based on  $w_{\max}$  within convective cores identified as regions with 20 dBZ or greater radar reflectivity. While the parameterization includes the high flash rate storms, thereby giving a frequency distribution shaped similar to that observed without the erroneous drop-off, the flash count is under-predicted by up to a factor of 10. From this experiment, we see the need to evaluate flash rate parameterizations with observations for the locations and periods specific to the simulations. It is insufficient to use high resolution model results as “truth” for coarse resolution simulations. Hence, validation and tuning prior to further usage of  $f(w_{\max})$  from Eq. (2) is encouraged. Furthermore, parameterizing flash rate in cloud-resolving models based on other storm parameters (Barthe et al., 2010) should also be tested.

To summarize, we recommend the following when applying the Prince and Rind parameterizations for lightning flash rates:

1. Proper adjustment to cloud top should be made to match the expected 20 dBZ radar reflectivity top when applying PR92;
2. PR93 for IC : CG partitioning should be used only if it is unnecessary to get information on the frequency distribution of flashes;
3. Scaling for resolution dependency may be performed by areal ratio against a base-case resolution, defined as that producing 1 storm per grid within the domain of interest (36 km grid spacing in this study).

To further the confidence of the lightning flash rate parameterizations and IC : CG partitioning, long-term wide-area total lightning detection and data archiving should be accompanied by coincident observations of cloud-top or other convective properties with well-defined error characteristics in observations and quantifiable predictability in numerical models.

## Appendix A

### Comments on cloud-top height reduction

In this study, we used the level of neutral buoyancy (LNB) from the WRF implementation of the Grell–Devenyi convective parameterization (Grell and Devenyi, 2002) as a proxy for sub-grid cloud-top heights for the purpose of testing a flash rate parameterization by Price and Rind (1992, 1993, 1994). A reduction of 2 km is used to reconcile the differences between LNB and the cloud top that would be obtained if defined at a 20 dBZ reflectivity threshold. While this

method produces an integrated flash count consistent with that observed after taking into account model biases in convective precipitation, we acknowledge that storm-to-storm variability cannot be captured by such a simple approach. Presented in this section are offline calculations of both 20 dBZ cloud tops and LNB cloud tops from a 13-day simulation at 4 km grid spacing to understand the margin of potential errors.

Radar reflectivity is estimated by using rain, snow, and graupel particle information from hourly outputs. For consistency, the offline calculation of reflectivity uses the same modified equations from Smith et al. (1975) and criteria as those used in the 4 km simulation. The highest model level with more than 20 dBZ is then defined as the 20 dBZ top.

LNB is estimated by a simple “parcel method,” rather than emulating the full algorithm in the parameterization as implemented in WRF. Therefore, the result may differ from what would be produced within the model. First, the dew point depression at the surface model level is determined, which is then used to seek the lifting condensation level (LCL) assuming adiabatic ascent. From the LCL, the moist adiabatic lapse rate ( $\Gamma_m$ ) is calculated and the level of free convection (LFC) is determined by linearly extrapolating the moist adiabat using the lower level’s  $\Gamma_m$  to the model level immediately above. From the LFC, a search is performed at incremental model levels until the LNB is exceeded. Grid points with LFCs < 500 m or above-freezing temperature at LNBs are discarded.

In total,  $1.34 \times 10^6$  columns with sufficient reflectivity and cloud-top heights greater than 5 km AGL are found. The distribution of the difference  $\langle h_{\text{LNB}} - h_{\text{dBZ}} \rangle$  indicates that LNB is higher than the 20 dBZ top 62 % of the time with a mean of 1.1 km and a standard deviation of 2.3 km. Other metrics for defining the required offsets between the two heights can produce different results. For example, to minimize the bias after applying PR92 with  $|\langle (h_{\text{LNB}} - \delta h)^5 / h_{\text{dBZ}}^5 \rangle - 1|$ , the reduction  $\delta h$  evaluates to 3.27 km. While the 2 km reduction used in this study differs from the two computed here, it is within the calculated range and thus can still be considered a median representation, especially when the uncertainties in the methods used for the offline LNB and radar reflectivity computations are taken into account.

Finally, it is essential to re-emphasize that the choice of cloud-top reduction is specific to the use of the Grell–Devenyi convective parameterization in WRF or other models producing LNB as the best-available proxy for sub-grid cloud-tops. In other models, cloud-top proxies other than LNB may be present. In those cases, an adjustment specific to those proxies should be used if PR92 is the preferred method for parameterization. An alternative would be to reformulate PR92 to be based on LNB, which lies beyond the scope of this paper and may be attempted in the future as needed.

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